FLUID PRESSURE AT A POINT

Consider a small area dA in large mass of fluid. If the fluid is stationary, then the force exerted by the surrounding fluid on the area dA will always be perpendicular to the surface dA. Let dF is the force acting on the area dA in the normal direction. Then the ratio of $\frac{dF}{dA}$ is known as the intensity of pressure or simply pressure and this ratio is represented by p. Hence mathematically the pressure at a point in a fluid at rest is

$$p = \frac{dF}{dA}.$$

If the force (F) is uniformly distributed over the area (A), then pressure at any point is given by

$$p = \frac{F}{A} = \frac{\text{Force}}{\text{Area}}$$

 \therefore Force or pressure force, $F = p \times A$.

The units of pressure are: (i) kgf/m² and kgf/cm² in MKS units, (ii) Newton/m² or N/m² and N/mm² in SI units. N/m² is known as Pascal and is represented by Pa. Other commonly used units of pressure are:

$$kPa = kilo pascal = 1000 N/m^2$$

 $bar = 100 kPa = 10^5 N/m^2$.

PASCAL'S LAW

It states that the pressure or intensity of pressure at a point in a static fluid is equal in all directions. This is proved as:

The fluid element is of very small dimensions i.e., dx, dy and ds.

Consider an arbitrary fluid element of wedge shape in a fluid mass at rest as shown in Fig. 2.1. Let the width of the element perpendicular to the plane of paper is unity and p_x .

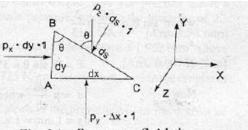


Fig. 2.1 Forces on a fluid element.

 p_y and p_z are the pressures or intensity of pressure acting on the face AB, AC and BC respectively. Let $\angle ABC = \theta$. Then the forces acting on the element are:

1. Pressure forces normal to the surfaces.

2. Weight of element in the vertical direction.

The forces on the faces are:

Force on the face
$$AB = p_x \times \text{Area of face } AB$$

$$= p_x \times dy \times 1$$

Similarly force on the face $AC = p_y \times dx \times 1$ Force on the face $BC = p_z \times ds \times 1$

Weight of element = (Mass of element) $\times g$

= (Volume
$$\times \rho$$
) $\times g = \left(\frac{AB \times AC}{2} \times 1\right) \times \rho \times g$,

where $\rho = density of fluid.$

or

Resolving the forces in x-direction, we have

$$p_x \times dy \times 1 - p (ds \times 1) \sin (90^\circ - \theta) = 0$$

or $p_x \times dy \times 1 - p_z ds \times 1 \cos \theta = 0$. But from Fig. 2.1, $ds \cos \theta = AB$

$$ds \cos \theta = AB = dy$$
$$p_x \times dy \times 1 - p_z \times dy \times 1 = 0$$

 $p_x = p_z$

Similarly, resolving the forces in y-direction, we get

$$p_y \times dx \times 1 - p_z \times ds \times 1 \cos(90^\circ - \theta) - \frac{dx \times dy}{2} \times 1 \times \rho \times g = 0$$

 $p_{y} \times dx - p_{z} ds \sin \theta - \frac{dxdy}{2} \times \rho \times g = 0.$

But $ds \sin \theta = dx$ and also the element is very small and hence weight is negligible. $p_y dx - p_z \times dx = 0$

$$p_y dx - p_z \times dx = 0$$

$$p_y = p_z$$
...(2.2)

From equations (2.1) and (2.2), we have

$$p_x = p_y = p_z$$
 ...(2.3)

The above equation shows that the pressure at any point in x, y and z directions is equal.

PRESSURE VARIATION IN A FLUID AT REST

The pressure at any point in a fluid at rest is obtained by the Hydrostatic Law which states that the rate of increase of pressure in a vertically downward direction must be equal to the specific weight of the fluid at that point. This is proved as:

Consider a small fluid element as shown in Fig. 2.2

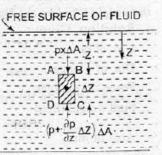
Let $\Delta A = \text{Cross-sectional}$ area of element

 ΔZ = Height of fluid element

p =Pressure on face AB

Z = Distance of fluid element from free surface.

The forces acting on the fluid element are



...(2.1)

Fig. 2.2 Forces on a fluid element.

- 1. Pressure force on $AB = p \times \Delta A$ and acting perpendicular to face AB in the downward direction.
- 2. Pressure force on $CD = \left(p + \frac{\partial p}{\partial Z}\Delta Z\right) \times \Delta A$, acting perpendicular to face CD, vertically upward direction.
 - 3. Weight of fluid element = Density $\times g \times \text{Volume} = \rho \times g \times (\Delta A \times \Delta Z)$.
- 4. Pressure forces on surfaces BC and AD are equal and opposite. For equilibrium of fluid element, we have

$$p\Delta A - \left(p + \frac{\partial p}{\partial Z}\Delta Z\right)\Delta A + \rho \times g \times (\Delta A \times \Delta Z) = 0$$
or
$$p\Delta A - p\Delta A - \frac{\partial p}{\partial Z}\Delta Z\Delta A + \rho \times g \times \Delta A \times Z = 0$$
or
$$-\frac{\partial p}{\partial Z}\Delta Z\Delta A + \rho \times g \times \Delta A\Delta Z = 0$$
or
$$\frac{\partial p}{\partial Z}\Delta Z\Delta A = \rho \times g \times \Delta A\Delta Z \text{ or } \frac{\partial p}{\partial Z} = \rho \times g \text{ [cancelling } \Delta A\Delta Z \text{ on both sides]}$$

$$\vdots \qquad \frac{\partial p}{\partial Z} = \rho \times g = w \qquad (\because p \times g = w) \qquad \dots (2.4)$$

where w = Weight density of fluid.

Equation (2.4) states that rate of increase of pressure in a vertical direction is equal to weight density of the fluid at that point. This is **Hydrostatic Law**.

By integrating the above equation (2.4) for liquids, we get

$$\int dp = \int \rho g Z$$

$$p = \rho g Z \qquad ...(2.5)$$
Althoughberic pressure and Z is the ball to S.d.

where p is the pressure above atmospheric pressure and Z is the height of the point from free surfaces.

From equation (2.5), we have
$$Z = \frac{p}{\rho \times g}$$
 ...(2.6)

Here Z is called pressure head.